

Real-time dynamics in Quantum Impurity Systems: A Time-dependent Numerical Renormalization Group Approach

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What is a Quantum-Impurity System (QIS)?

Quantum Impurity

== $|\alpha\rangle$

== $|\gamma\rangle$

(metallic) host
bosonic bath

quantum-impurity:

- embedded in a (metallic) host
- interacting with the environment of non-interacting particles (Bosons/Fermions)

Problem:

- infrared divergence due to local degeneracy

What is a Quantum-Impurity System (QIS)?

Quantum Impurity

≡ $|\alpha\rangle$

≡ $|\gamma\rangle$

(metallic) host
bosonic bath

Examples:

- transition metal ion Cu, Mn, Ce in a metal
- two-level system (Qubit) in a bosonic bath
- Quantum dot coupled to leads
- donor-acceptor centers of a large bio-molecule
- ...

Goal of the Talk

Our new Approach to Non-Equilibrium of QIS:

- based on the **non-perturbative NRG**
 - uses the ***complete basis of the many body Fock space***
 - takes into account ***all energy scales***
 - describes **short and long** time scales
 - does **not accumulate an error $\propto t$** as the TD-DMRG
- ➡ **breakthrough in the description of real time dynamics of non-equilibrium quantum systems:**

Contents

1. Introduction

- Modelling of quantum dots
- Charge transfer in molecules (spin-boson model)

2. Non-equilibrium dynamics

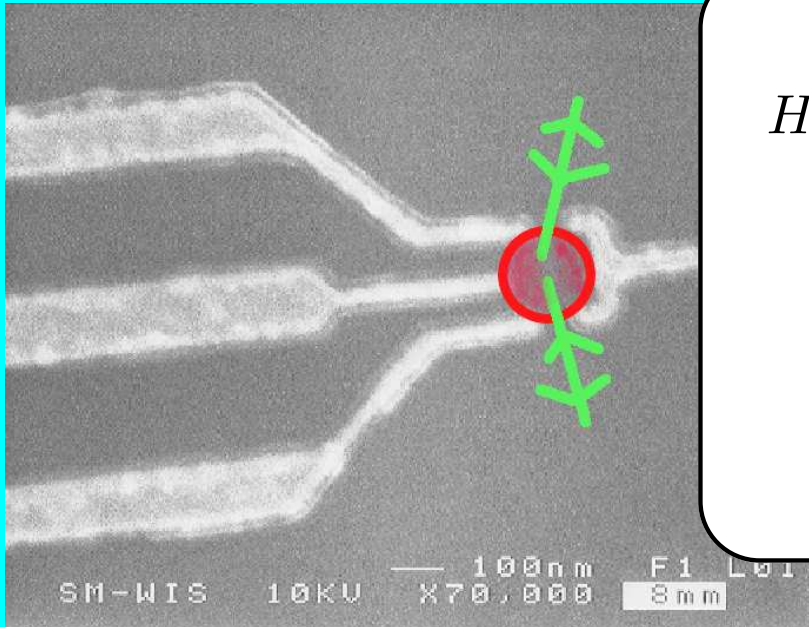
- Time evolution of quantum systems
- New approach to quantum impurity problems

3. Results

- Dissipation and decoherence in a two level system
- Spin- and charge dynamics in ultra-small quantum dots
- AF-Kondo model
 - spin precession

4. Summary and outlook

Modelling of a Quantum Dot



$$H = \sum_{k\sigma} \epsilon_{k\sigma} c_{k\sigma}^\dagger c_{k\sigma} + \sum_{\sigma} [E_d(t) - \sigma H(t)] d_{\sigma}^\dagger d_{\sigma} + U n_{\uparrow}^d n_{\downarrow}^d + \sum_{k\sigma} V(t) (c_{k\sigma}^\dagger d_{\sigma} + d_{\sigma}^\dagger c_{k\sigma})$$

Single Impurity Anderson Model (SIAM)

- charge fluctuation scale: $\Gamma_i = V_i^2 \pi \rho_F$
- infrared problem
 - ☞ low temperature scale: $T_K \propto \exp(-\pi U/8\Gamma)$

Spin-Boson Model

- qubit plus environment (Unruh)
- electron transfer in (bio)-molecules (Marcus, Schulten)

$$|\uparrow\rangle = |A\rangle, |\downarrow\rangle = |D\rangle$$

$$H = \epsilon\sigma_z - \frac{\Delta}{2}\sigma_x + \sum_q \hbar\omega_q b_q^\dagger b_q + \sigma_z \sum_q M_q (b_q^\dagger + b_q)$$

$$J(\omega) = \sum_q |M_q|^2 \delta(\omega - \omega_q) \propto 2\pi\alpha\omega_c^{1-s}\omega^s$$

Leggett *et. al.* (RMP 1987), Xu and Schulten 1994, Bulla *et. al.* (2003) ...

Questions:

- influence of the bosonic spectrum $J(\omega)$ on the real time dynamics
- critical slowdown of the charge transfer process for large coupling

Where do we stand in the
description of non-equilibrium,
dissipation and decoherence
in quantum systems?

Non-Equilibrium Dynamics of Quantum Systems

quantum dynamics

- single quantum state: Schrödinger equation

$$i\hbar\partial_t|\psi\rangle = H(t)|\psi\rangle$$

- ensemble: density operator

$$i\hbar\partial_t\hat{\rho}(t) = [H, \hat{\rho}] \quad ; \quad \rho(t) = e^{-iHt/\hbar}\rho_0e^{iHt/\hbar}$$

finite size quantum system: only unitary dynamics, no dissipation

dissipation and decoherence:
infinitely large environment needed

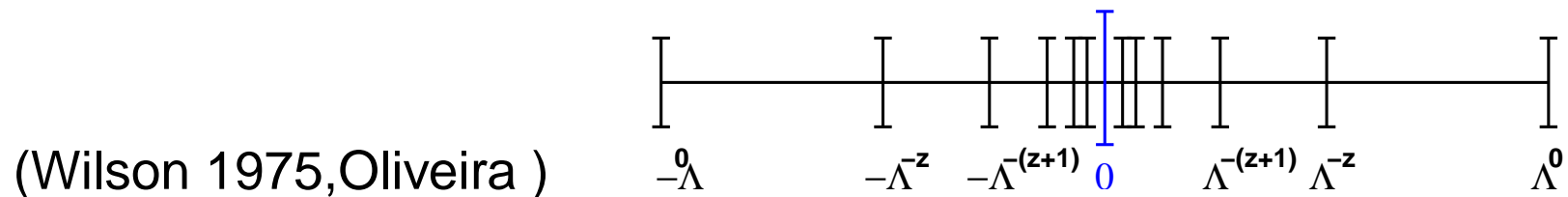
$$\frac{\text{Size of Subsystem}}{\text{Size of environment}} \longrightarrow 0$$



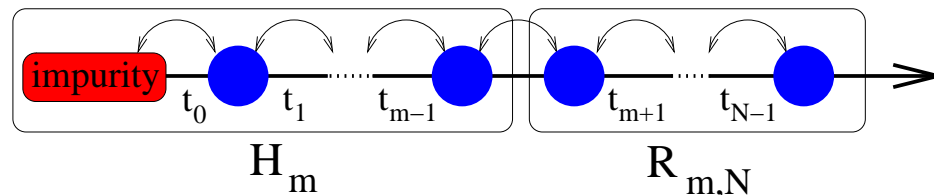
NRG Approach to Quantum Impurity Problems

$$H = H_{imp} + H_{bath} + H_{imp-bath}$$

1. discretizing the bath Hamiltonian on a logarithmic energy mesh



2. mapping onto a semi-infinite chain

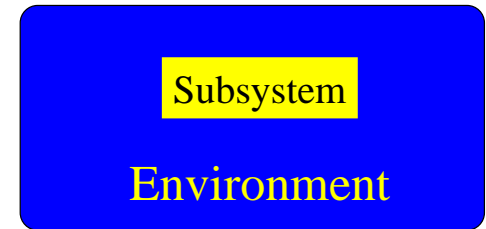
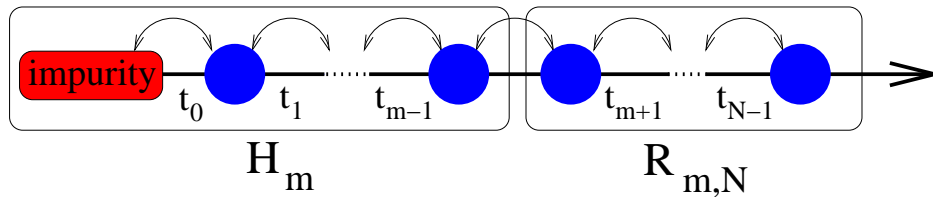


3. diagonalizing the Hamiltonian H_{N+1} using the recursion

$$H_{N+1} = \sqrt{\Lambda} H_N + \sum_{\alpha} \xi_{N\alpha} \left(f_{N+1\alpha}^{\dagger} f_{N\alpha} + f_{N\alpha}^{\dagger} f_{N+1\alpha} \right)$$

4. truncate the basis set, go back to step 3

Novel Many-Body Approach to NEQ of QIS



- use the NRG to generate a complete basis

$|l, e; m\rangle$

- $H_m |l\rangle = E_l^m |l\rangle$, l eliminated state

- $e \in R_{m,N}$

$$1 = \sum_m \sum_{l,e} |l, e; m\rangle \langle l, e; m|$$

- Puls at $t = 0$: $H(t) = H_i \Theta(-t) + H_f \Theta(t)$

- operator \hat{O} : property of the subsystem S

Novel Many-Body Approach to NEQ of QIS

time-dependent NRG (TD-NRG)

(FBA, A. Schiller, cond-mat/0505553, PRL 2005)

● calculate $\rho_{NEQ}^{red}(t)$

Subsystem

Environment

$$\langle \hat{O} \rangle(t) = \text{Tr} [\rho(t) \hat{O}] = \sum_{m, \alpha \alpha'} \langle \alpha | \hat{O} | \alpha' \rangle \rho_{\alpha \alpha', m}^{red}(t)$$

$$\rho_{\alpha \alpha', m}^{red}(t) = e^{-i(E_\alpha - E_{\alpha'})t} \sum_e \langle \alpha, e; m | \rho_{eq} | \alpha', e; m \rangle$$

Feynman 1972, White 1992, Hofstetter 2000, . . .

● mimic bath continuum: use Oliveira's z -trick

● evolves towards the new steady state: $[H(t > 0), \rho(\infty)] = 0$

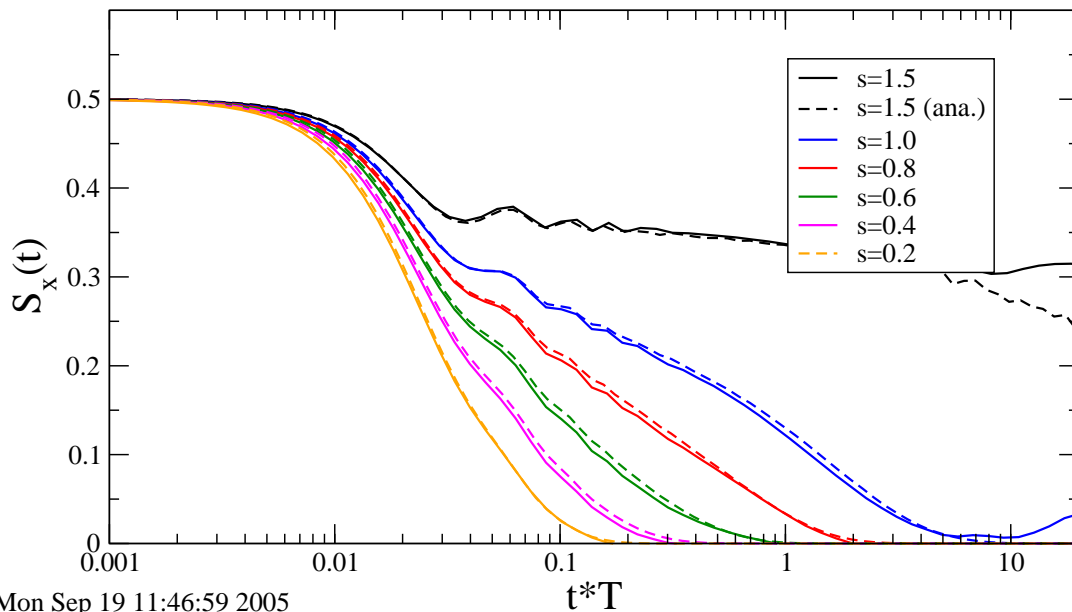
Trace over the environment: dissipation and decoherence!

Spin-Boson Model

$$H = \epsilon \sigma_z - \frac{\Delta}{2} \sigma_x + \sum_q \hbar \omega_q b_q^\dagger b_q + \sigma_z \sum_q M_q (b_q^\dagger + b_q)$$

$$S_x = \frac{1}{2} (|\uparrow\rangle\langle\downarrow| + |\downarrow\rangle\langle\uparrow|)$$

$N_s=150, N_z=16, N_b=8, N_{iter}=14, T=0.0078125, \Lambda=2^{1/2}, \alpha_{damp}=0.1, \alpha=0.1, \epsilon=0, \Delta^0=0., \omega_c=1$



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Decoherence



QuBit state

$$\frac{1}{\sqrt{2}} (|\uparrow\rangle + |\downarrow\rangle)$$



exact solution

$$P(t) = e^{-\Gamma(t)}$$

Leggett *et al.*, Unruh,
Palma *et al.*, Bulla *et al.*

Spin-Boson Model

$$H = \epsilon \sigma_z - \frac{\Delta}{2} \sigma_x + \sum_q \hbar \omega_q b_q^\dagger b_q + \sigma_z \sum_q M_q (b_q^\dagger + b_q)$$

$$J(\omega) = 2\pi \alpha \omega_c^{1-s} \omega^s \quad \text{for } 0 < \omega < \omega_c ;$$

Ohmic case: $s = 1$

Fixed point: delocalized

$$0 < \alpha < 1/2$$

oszillatory

Toulouse Point: $\alpha = 1/2$

$$1/2 < \alpha < \alpha_c(\Delta)$$

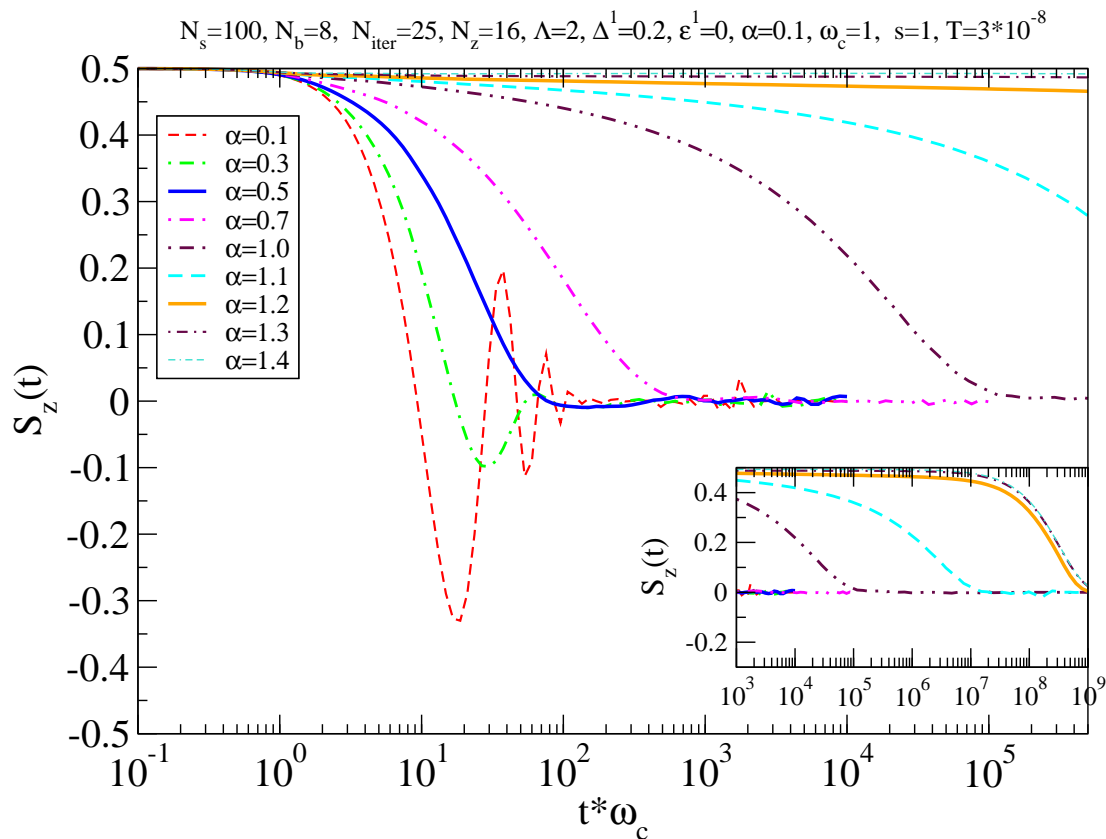
overdamped

localized

$$\alpha_c(\Delta) < \alpha$$

Spin-Boson Model

$$H = \epsilon \sigma_z - \frac{\Delta}{2} \sigma_x + \sum_q \hbar \omega_q b_q^\dagger b_q + \sigma_z \sum_q M_q (b_q^\dagger + b_q)$$

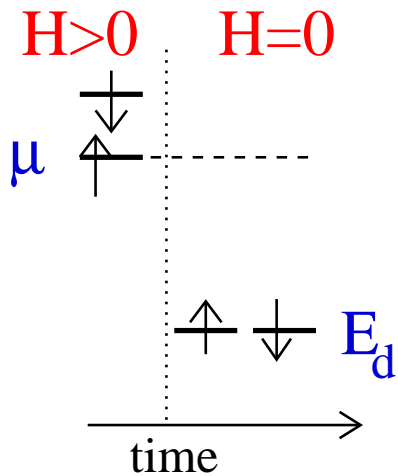


$$J(\omega) = 2\pi\alpha\omega_c^{1-s}\omega^s$$

Ohmic Regime: $s = 1$

- QPT at $\alpha_c(\Delta)$
- Toulouse point $\alpha = 1/2$
- oscillatory $\alpha < 1/2$
overdamped $\alpha_c > \alpha > 1/2$
- localize: $\alpha > \alpha_c$

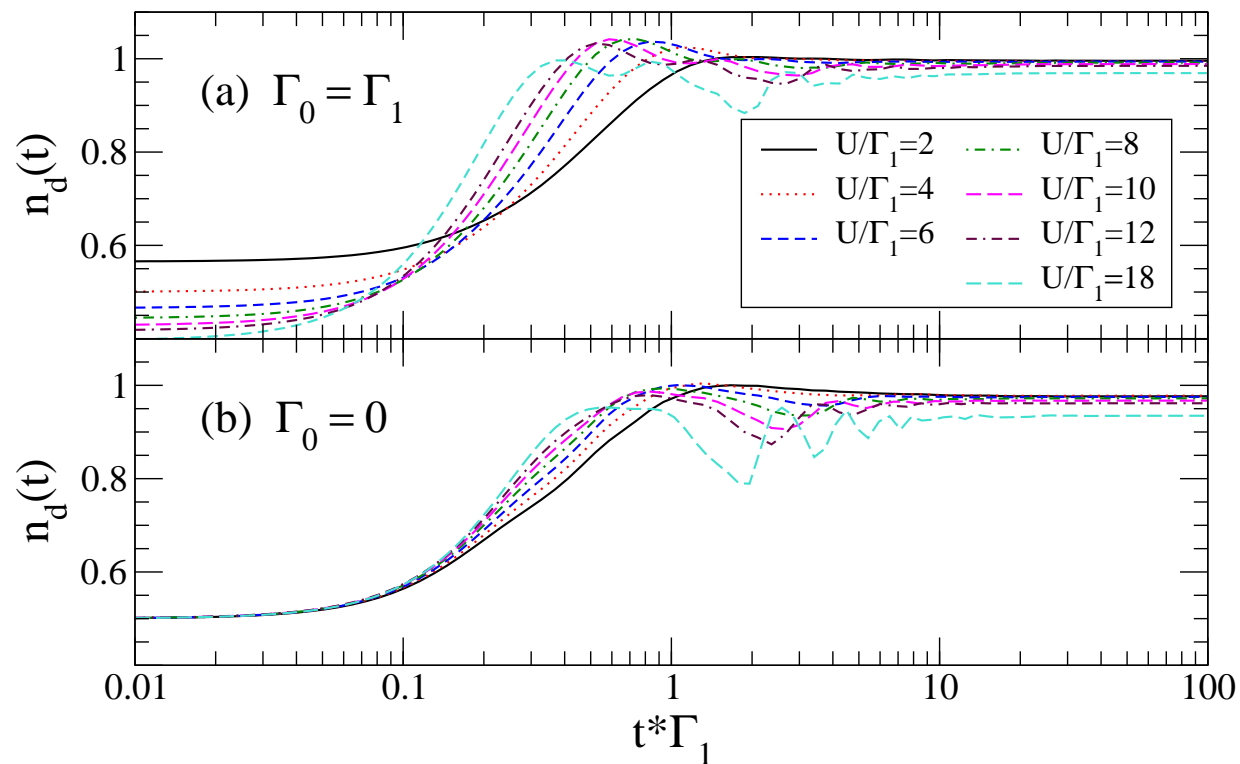
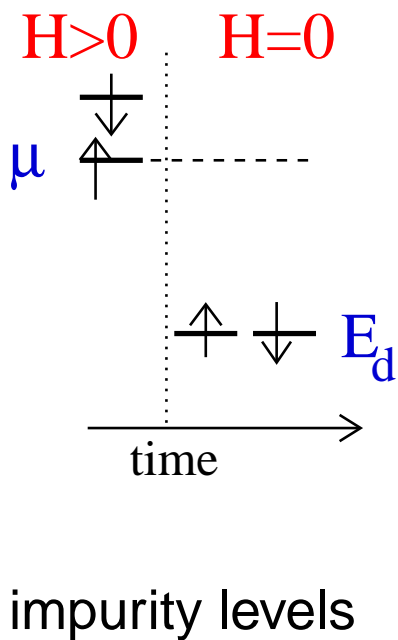
Charge Fluctuation in a Small Quantum Dot



$$\begin{aligned}
 H = & \sum_{k\sigma} \epsilon_{k\sigma} c_{k\sigma}^\dagger c_{k\sigma} \\
 & + \sum_{\sigma} [E_d(t) - \sigma H(t)] d_{\sigma}^\dagger d_{\sigma} + U n_{\uparrow}^d n_{\downarrow}^d \\
 & + \sum_{k\sigma} V(t) \left(c_{k\sigma}^\dagger d_{\sigma} + d_{\sigma}^\dagger c_{k\sigma} \right)
 \end{aligned}$$

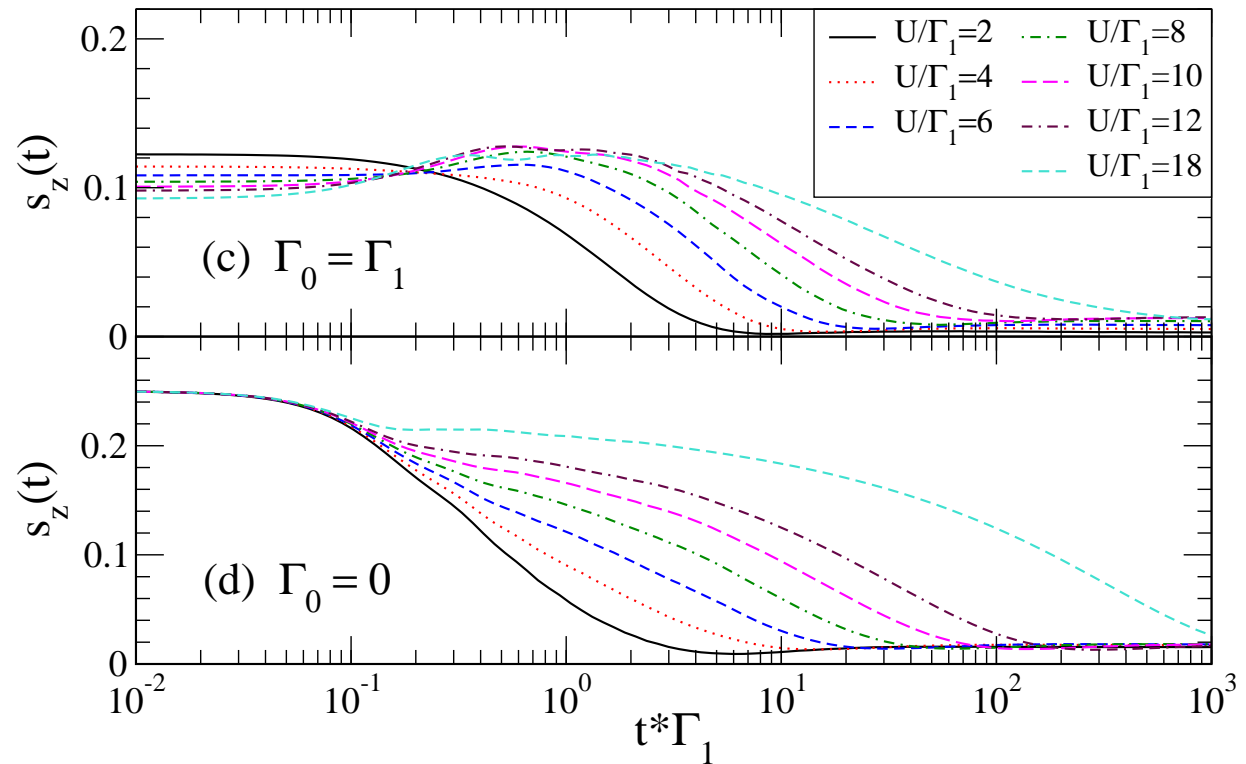
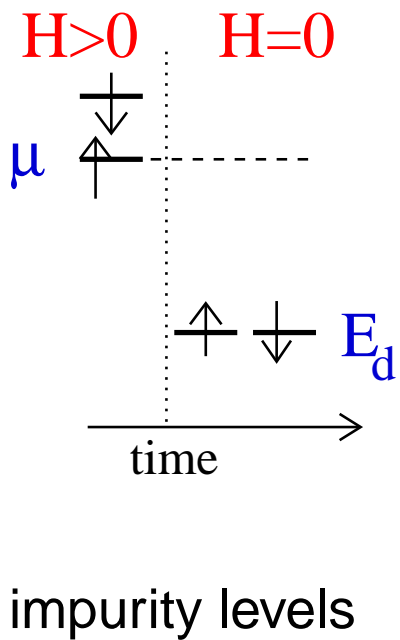
- change of E_d : change dynamics
- change of mag. field H : spin dynamics
- change of V : route to new equilibrium

Charge Fluctuation in a Small Quantum Dot

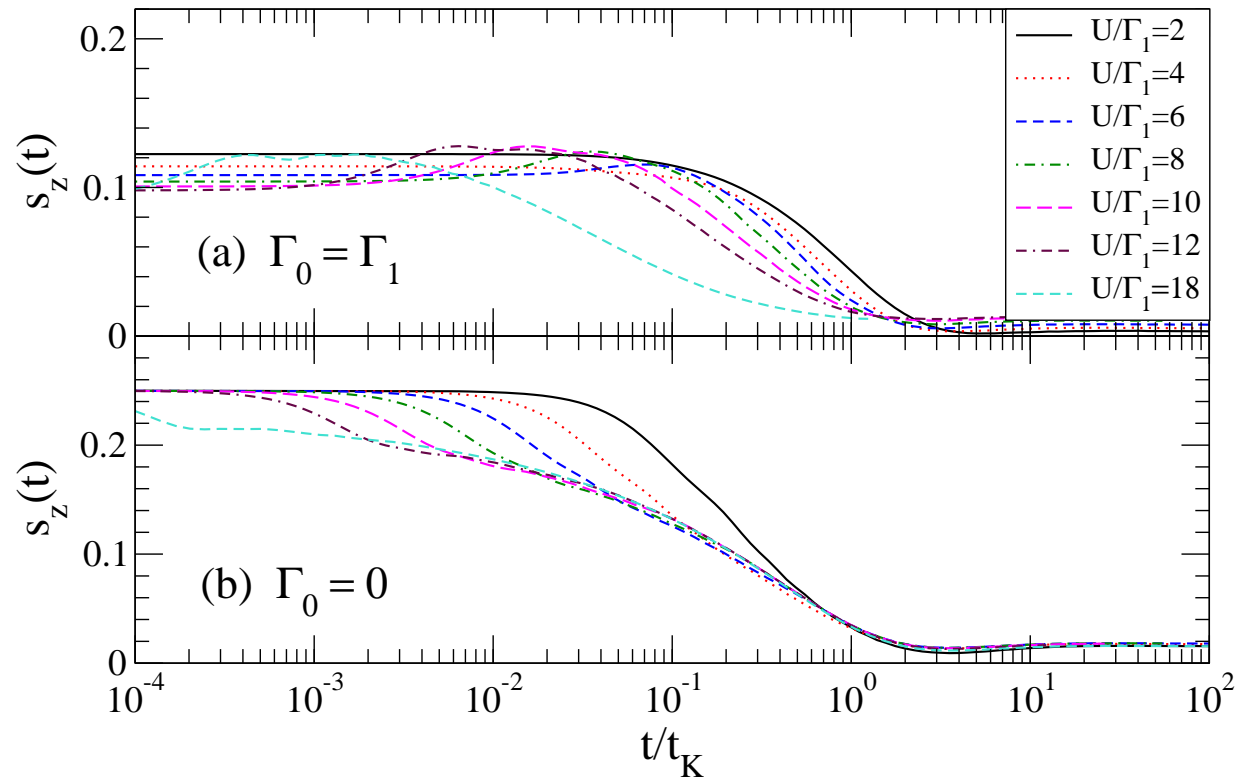
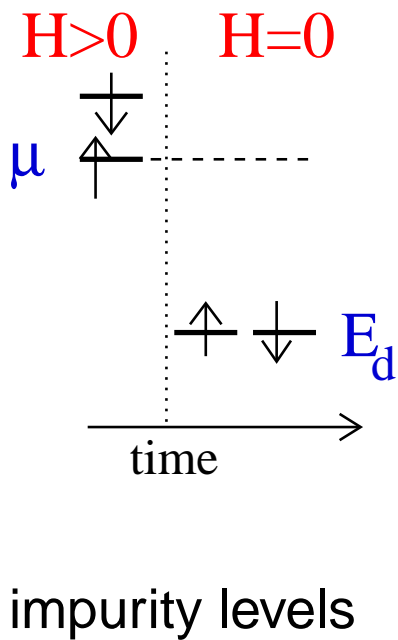


Charge relaxation time scale : $t_{ch} = 1/\Gamma_1$

Spin Fluctuation in a Small Quantum Dot

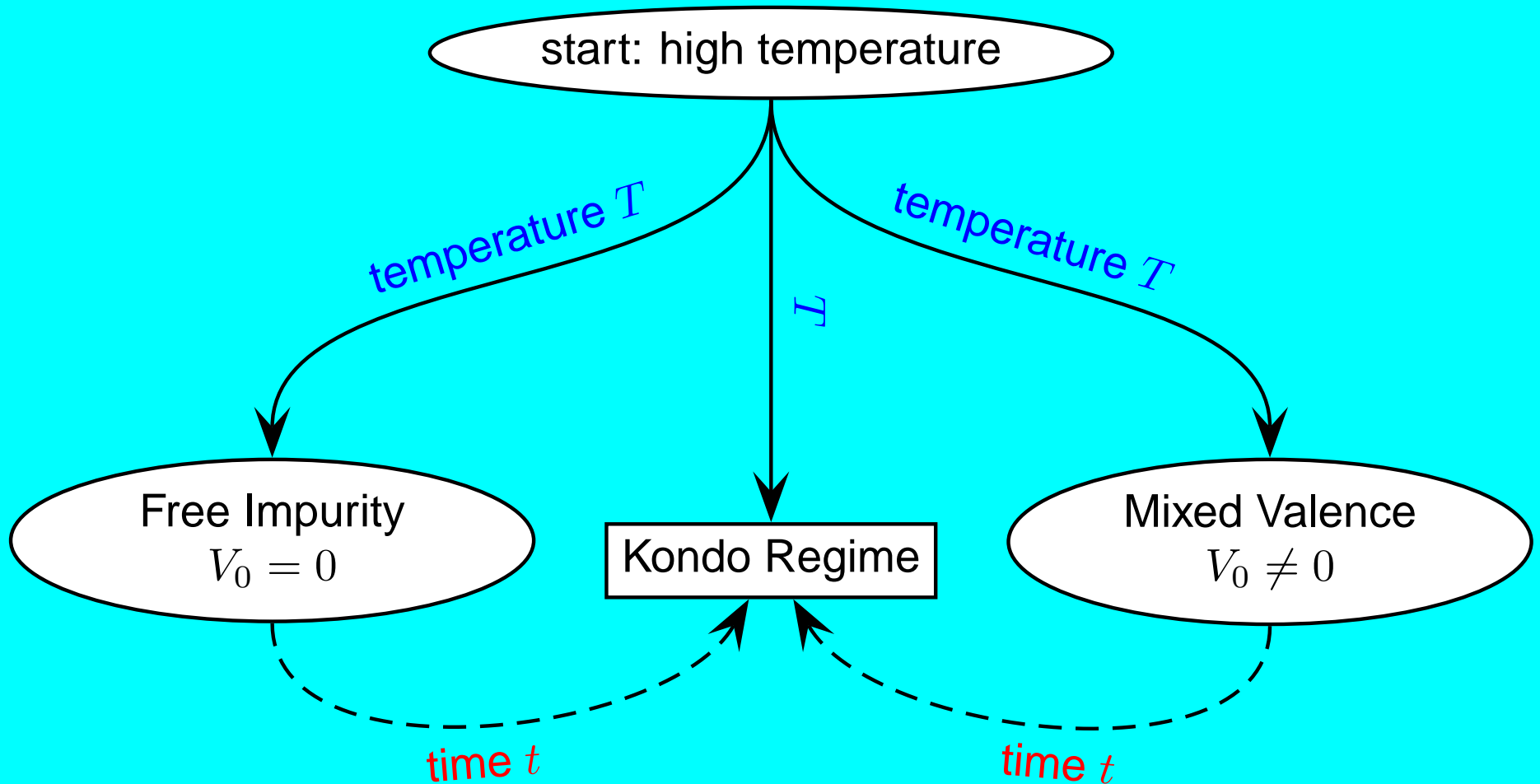


Spin Fluctuation in a Small Quantum Dot



Spin relaxation time scale : $t_{sp} \propto 1/T_K$

SIAM Time-Evolution

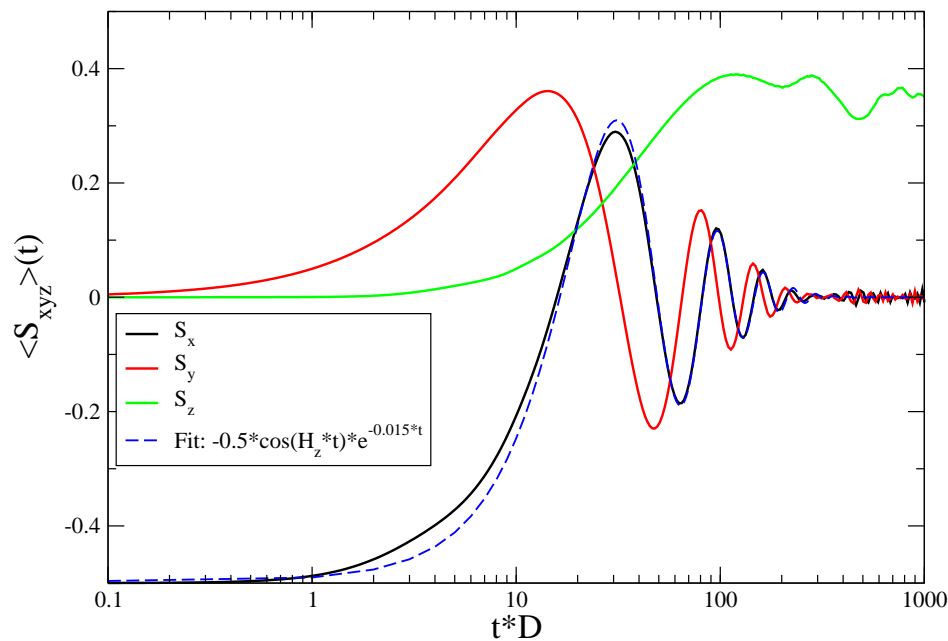


- no universality
- time evolution depends on boundary conditions

Kondo Model

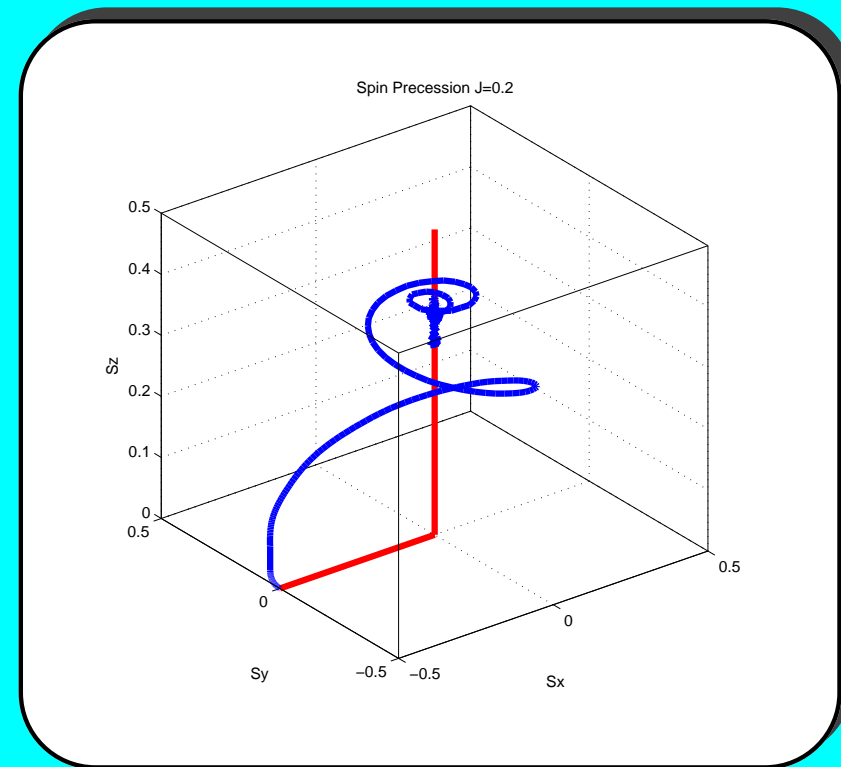
$$H = H_{cb} + \frac{H_z(t)}{2} \sigma_z + \frac{H_x(t)}{2} \sigma_x + J \vec{S}_{loc} \cdot \vec{s}_{cb}$$

Spin Precession: $H_x = 0.1 \rightarrow H_z$



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external magnetic field: $x \rightarrow z$ -axis



Summary and Outlook

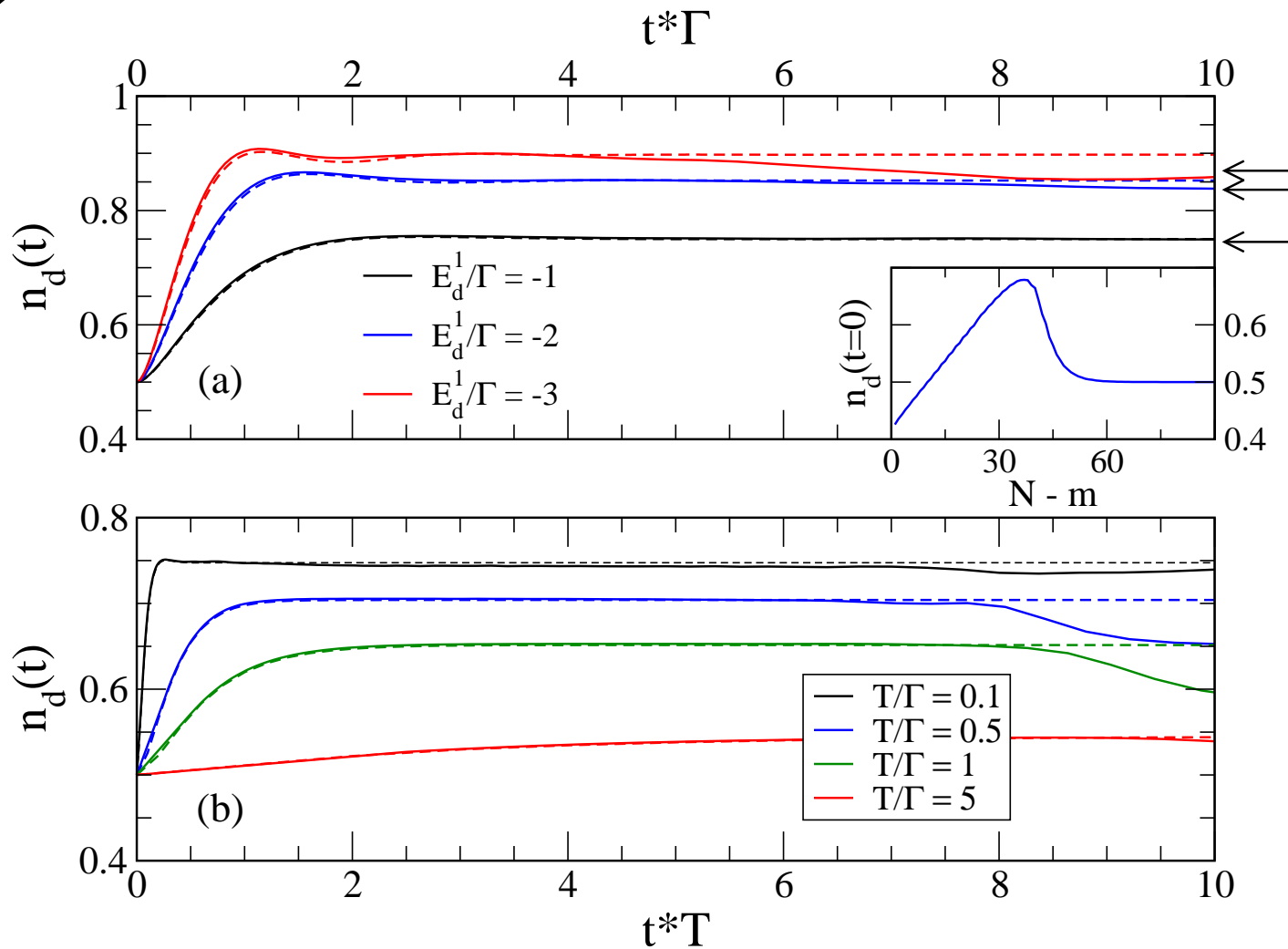
TD-NRG algorithm:

- **complete basis set** needed for the time evolution
- **dissipation and decoherence due to the bath**
- Spin-Boson Model
 - **benchmark:** decoherence and non-ohmic baths
 - oscillatory vs overdamped regime
- SIAM:
 - two relaxation time scale for spin and charge dynamics
- Kondo Model:
 - spin relaxation and spin precession

Outlook

- Bosonic QI models: charge transfer in bio-molecules, photosynthesis
- description of steady state currents

Benchmark: Resonant Level Model

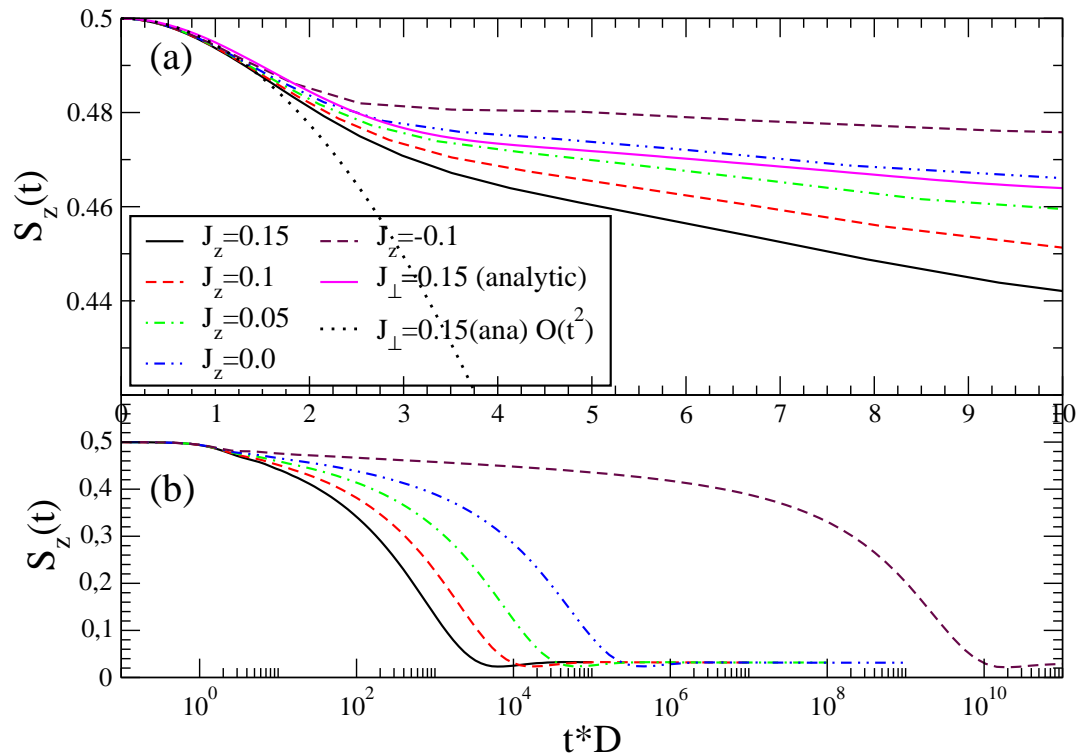


Anisotropic Kondo Model

$$H = H_{cb} + \frac{H_z(t)}{2} \sigma_z + \frac{H_x(t)}{2} \sigma_x + \sum_{kk' \alpha \beta} \frac{J_z(t)}{2} c_{k\alpha}^\dagger c_{k'\beta}^\dagger \sigma_z S_{imp}^z$$
$$+ \sum_{kk'} \frac{J_{perp}(t)}{2} \left(c_{k\uparrow}^\dagger c_{k'\downarrow}^\dagger S_{imp}^- + c_{k\downarrow}^\dagger c_{k'\uparrow}^\dagger S_{imp}^+ \right)$$

Anisotropic Kondo Model

spin relaxation: $J^0 = 0, J_{\perp}^1 = 0.15D = \text{const}$



two regimes:

● short time scale
analytical result:

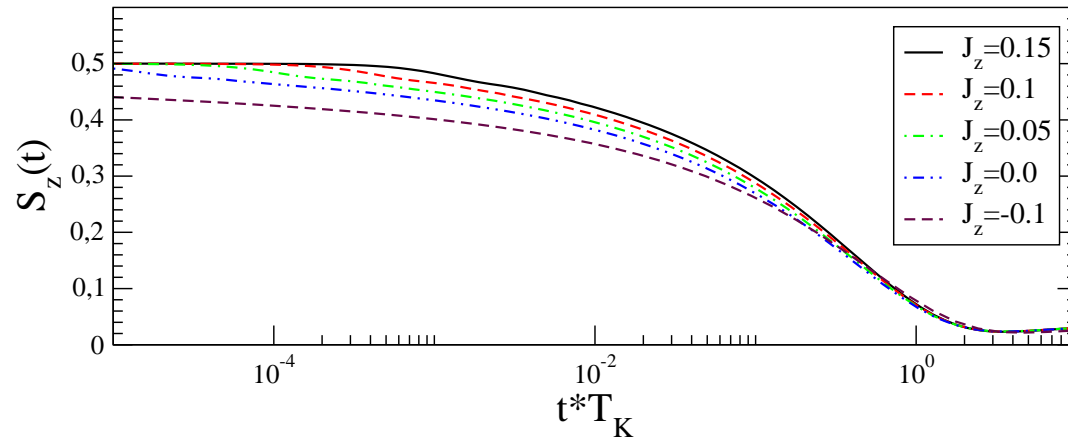
$$S_z(t) - S_z(0) \propto (2\rho J_{\perp}^1)^2 \times [G(2Dt) - 2G(Dt)]$$

$$G(x) = \sum_{l=1}^{\infty} \frac{(-1)^{l+1}}{(2l)!2l(2l-1)} x^{2l}$$

external magnetic field: z -axis

Anisotropic Kondo Model

spin relaxation: $J^0 = 0, J_{\perp}^1 = 0.15D = \text{const}$



two regimes:

● **short** time scale
analytical result:

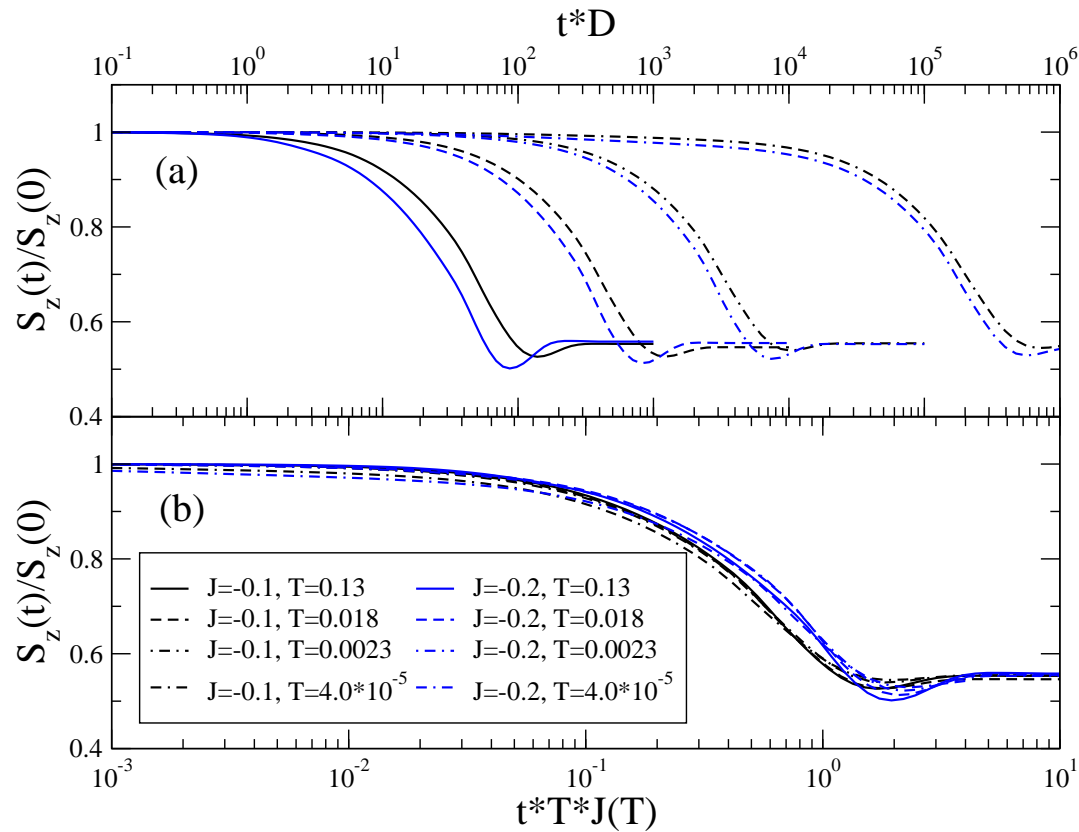
$$S_z(t) - S_z(0) \propto (2\rho J_{\perp}^1)^2 \times [G(2Dt) - 2G(Dt)]$$

$$G(x) = \sum_{l=1}^{\infty} \frac{(-1)^{l+1}}{(2l)!2l(2l-1)} x^{2l}$$

● **long** time:
 $t_{\text{long}} \propto 1/T_K$

Anisotropic Kondo Model

ferromagnetic regime, spin relaxation: $H_z = 0.01 \rightarrow 0$



external magnetic field: z -axis